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14. ABSTRACT In this study, we have employed a 2.5D time-dependent streamer-flux rope MHD model to investigate the role of magnetic reconnection in the CME acceleration processes. To implement this simulation, two distinct types of initial magnetic topology (i.e., Inverse and Normal prominence configuration) are used. The results showed: The normal prominence magnetic configuration forms a current sheet in the region of leading edge of the flux-rope, thus the magnetic reconnection occurs to open the closed field of the streamer which allows the flux rope to escape to launch a fast CME as suggested by Low and Zhang (2002). On the other hand, when the initial magnetic topology in the initial prominence configuration, the magnetic reconnection occurs at trailing edge of the flux-rope. In this case, the streamer confined the flux-rope until the magnetic reconnection occurs, this confined force removed which enlaced the CME acceleration (Wu et al. 2004) Both inverse and normal prominence configurations are able to produce fast CMEs.					
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AAMU Physics-based Analysis Using HASDM  
Agreement Dated 9/18/02  
UAH Account # 5-22187

*Final*

By

*S. T. Wu*

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During this period of performance June 15, 2002 – June 14, 2005, we have studied the role of magnetic reconnection in CME acceleration. The preliminary results were presented as an invited paper at the WISE (World Institute of Space Environment) Workshop on Computing in Space and Astrophysical Plasmas, April 18-22, 2005 in Katholieke Universiteit, Leuven, Belgium. A copy of the presented is attached.

# The Role of Magnetic Reconnection in CME Acceleration

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## Background

The analysis of LASCO/SOHO, Solar Maximum Mission (SMM) and Skylab observations show that there are many CMEs initiated with streamer and flux-rope magnetic topology (Dere et al. 1999; St. Cyr et al., 1999; Plunkett et al., 2000). Two types of CMEs have been distinguished with different kinematic characteristics (MacQueen and Fisher, 1983; Andrews and Howard, 2001). These are fast CMEs with high initial speeds (i.e. constant speed) and slow CMEs with low initial speeds but gradual acceleration (i.e. accelerated CMEs).

Efforts have been made to probe the underlying physics responsible for these dual characteristics. Low and Zhang (2002) proposed that fast and slow CMEs result from initial topology of the magnetic field characterized by normal and inverse quiescent prominences, respectively. They further suggested that when the flux-rope expulsion in the normal configuration of the prominence, the fast CME results and otherwise, the slow CME is launched. Liu et al (2003) have performed a numerical simulation to support this scenario. Wu et al. (2004) have performed numerical MHD study with various boundary conditions which mimic the solar surface conditions, these numerical experiments showed that flux-rope expulsion in the inverse configuration of the prominence could launch both types of CMEs (i.e. fast and slow).

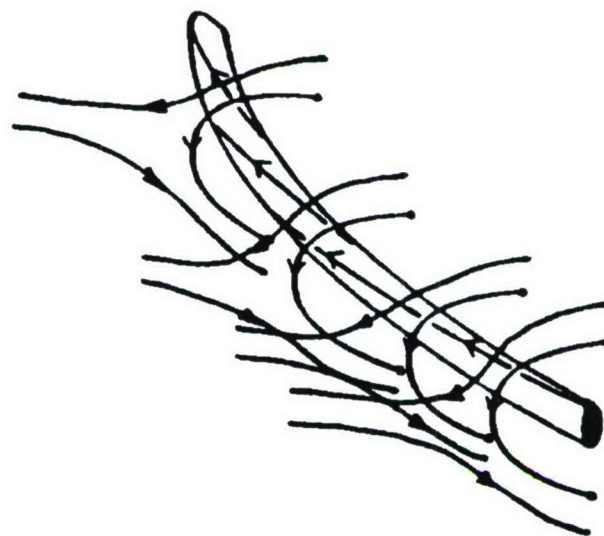
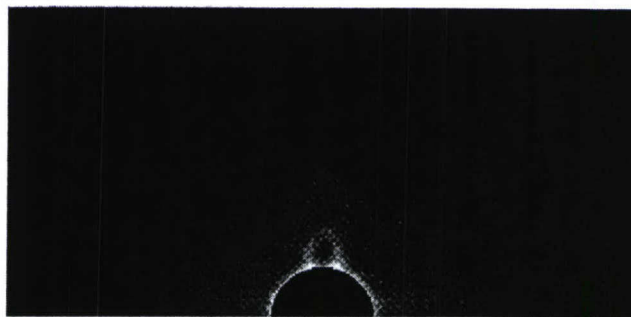
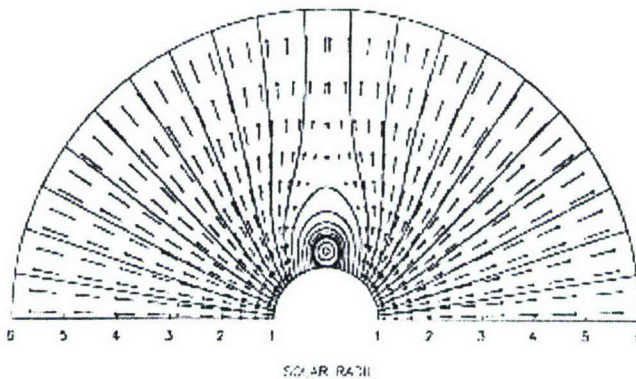
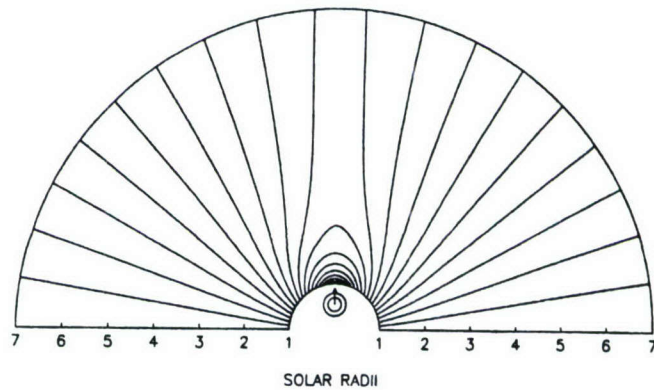
In the present study, we will show the role of magnetic reconnection in the CME acceleration processes occurs in these two types of initial topology of the magnetic field characterized by normal and inverse quiescent prominence, respectively.



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Initial State

## Initial process



Simulated initial state dynamical equilibrium state for a streamer and flux-rope system; (a) magnetic configuration of the streamer and flux-rope to be merged in (b) the dynamical equilibrium state of the magnetic field lines and velocity vector of the streamer, flux-rope and solar wind, (c) the polarization brightness of streamer and flux-rope structure and (d) the schematic description of the system.



## Properties of Initial State

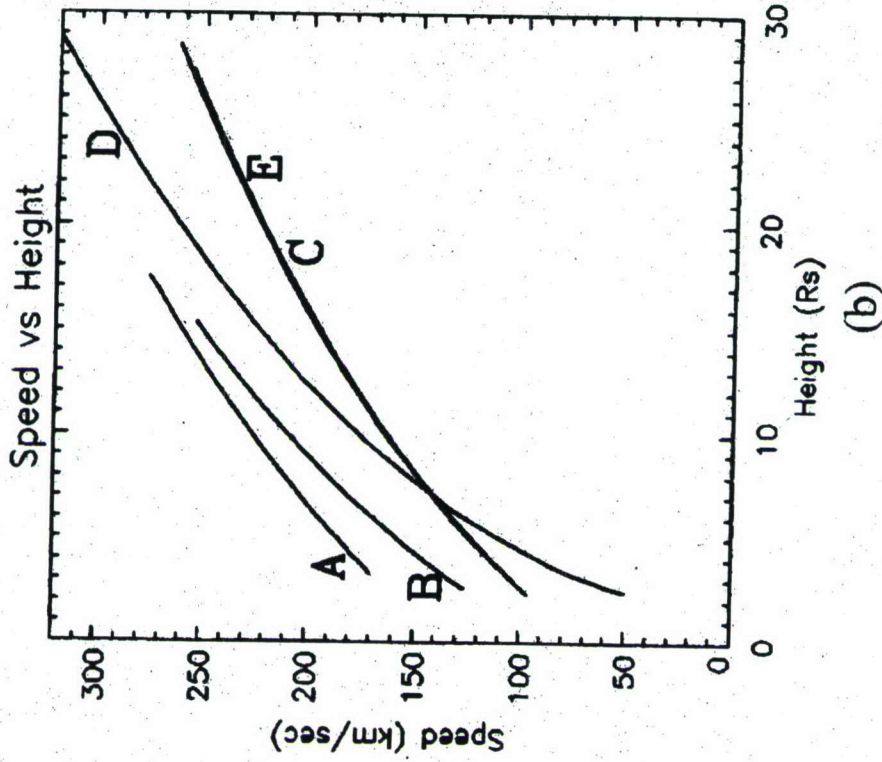
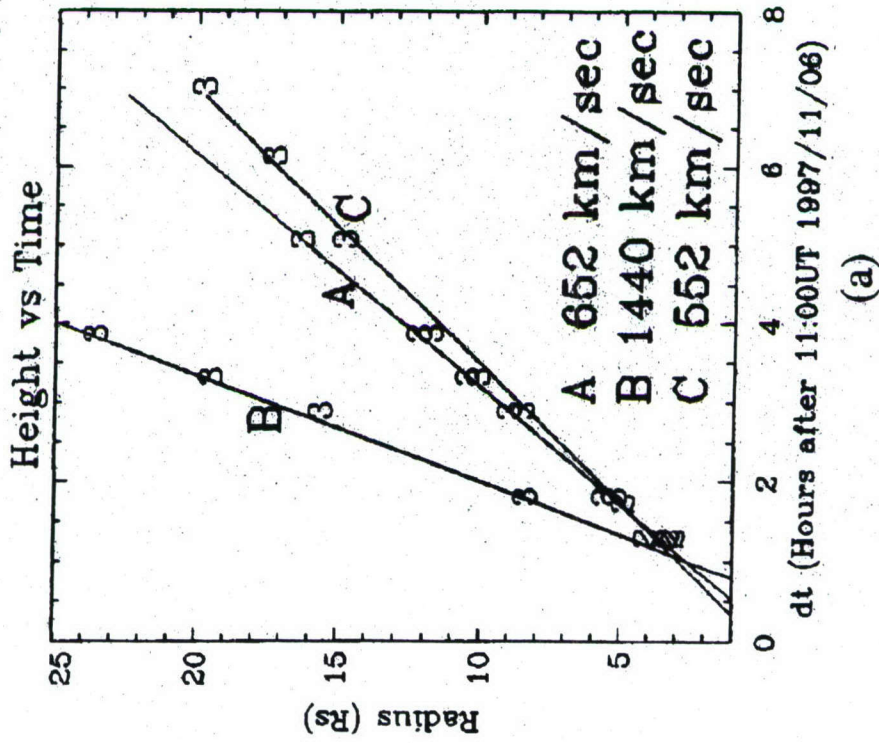
**Table Ia** Initial plasma properties, magnetic field strength, characteristic speeds at the equator and pole for the streamer and flux-rope system.

	Pole		Equator	
	1 R <sub>s</sub>	32 R <sub>s</sub>	1 R <sub>s</sub>	32 R <sub>s</sub>
N <sub>0</sub> (cm <sup>-3</sup> )	10 <sup>8</sup>	3.8 × 10 <sup>2</sup>	10 <sup>8</sup>	3.78 × 10 <sup>2</sup>
T <sub>0</sub> (MK)	1.4	0.77	1.4	0.768
B <sub>0</sub> (Gauss)	2.8	6.4 × 10 <sup>-4</sup>	1.14	5.62 × 10 <sup>-4</sup>
β	0.06	2.5	0.37	3.18
V <sub>sw</sub> (km/s)	108	80	108	79.6
V <sub>A</sub> (km/s)	609	72	249	63.1
V <sub>sw</sub> (km/s)	3.58	360	0	323
C <sub>f</sub> (km/s)	609	80	270	80
C <sub>s</sub> (km/s)	108	72	28	63

**Table Ib** Initial mass and energies for the streamer and flux-rope system.  
The energy is computed based on the depth being 0.1 R<sub>s</sub>.

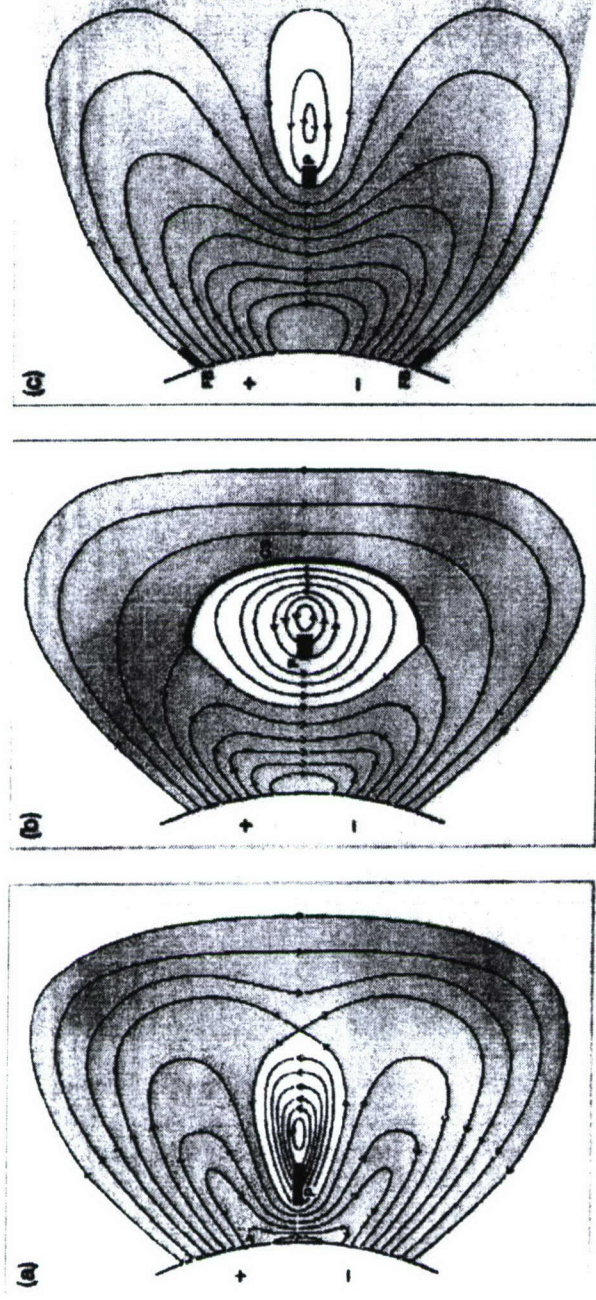
	Energy (10 <sup>30</sup> Ergs)			Mass (10 <sup>16</sup> g)	Total Energy (10 <sup>30</sup> ergs)
	Magnetic	Thermal	Kinetic		
Streamer	3.66	0.77	0.60	0.52	8.45
Rope	2.0	0.27	0	0.17	3.81
Rope + Streamer	5.66	1.04	0.60	0.69	12.26

## SOHO/LASCO Observation



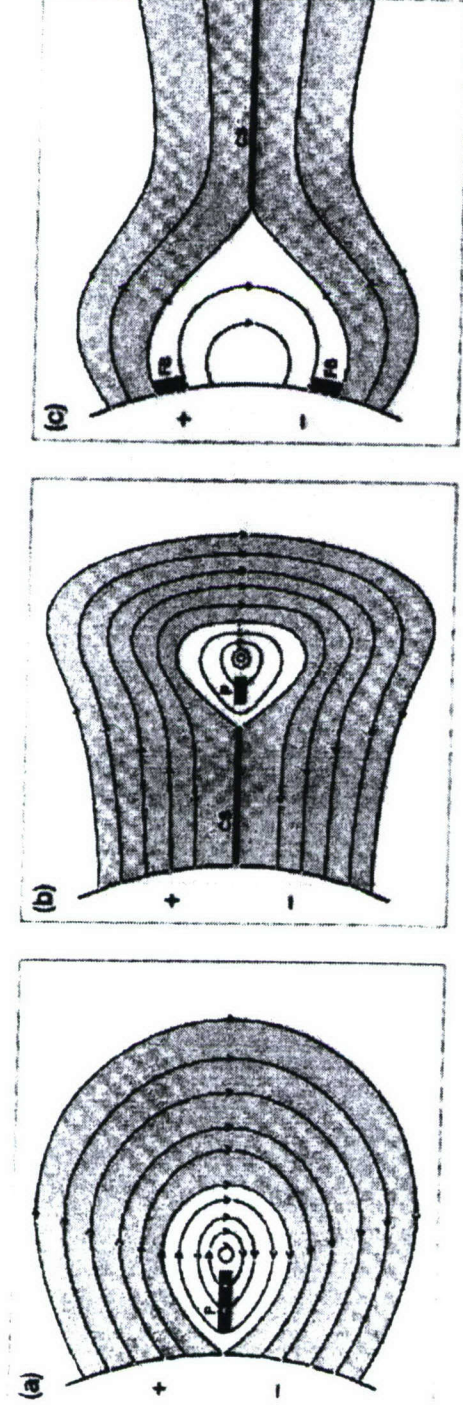
(a) The height-time curves for constant speed CME and (b) the velocity-distance curves for accelerated CMEs (Andrews and Howard, 2000).





Flux rope expulsion in the **normal configuration** (Low and Zhang, 2002) **(a)** The initial state prior to eruption; **(b)** the lift-off of the flux rope owing to mass-loss to the prominence forming a current sheet ahead of the flux rope and **(c)** further development with magnetic reconnection having removed the fields ahead of the flux rope to create the sling-shot field topology; a pair of footpoints brightenings (labeled FB) belonging to newly reconnected flux corresponds to the two-ribbon flare. Courtesy Low and Zhang.



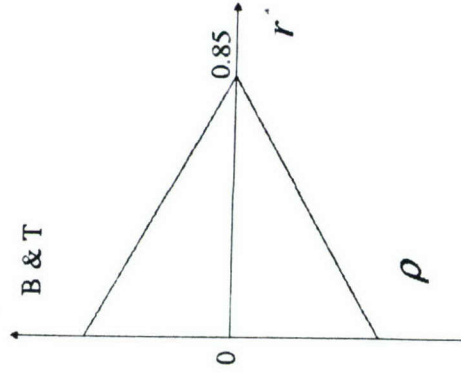


Flux rope eruption in the **inverse configuration** (Low and Zhang, 2002) **(a)** The initial state prior to eruption with a massive prominence sheet in the flux rope surrounded by dense coronal mass outside the flux rope; **(b)** the lift-off of the flux rope owing to loss of mass to the prominence forming a current sheet behind the flux rope and **(c)** further development after the flux rope expulsion, with magnetic reconnection producing closed bipolar fields anchored to the atmospheric base, showing a pair of footpoint brightenings (labeled FB) belonging to newly reconnected flux to correspond to the two-ribbon flare. Courtesy Low and Zhang.

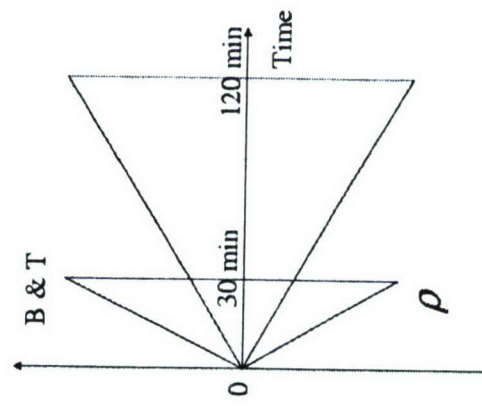
## Initiation Processes

To initiate this simulation, we will inject addition magnetic flux into the flux-rope and/or drain the plasma from the flux-rope. The numerical procedure to inject magnetic flux and drain the plasma are:

Spatial



Temporal



Test of the rate of  
the input

$$B_{\varphi}^{n+1} = B_{\varphi}^n \left[ 1 + \delta_B \left( 1 - \frac{r^*}{0.85r_f} \right) \right]$$

$$\rho^{n+1} = \rho^n \left[ 1 - \delta_{\rho} \left( 1 - \frac{r^*}{0.85r_f} \right) \right]$$

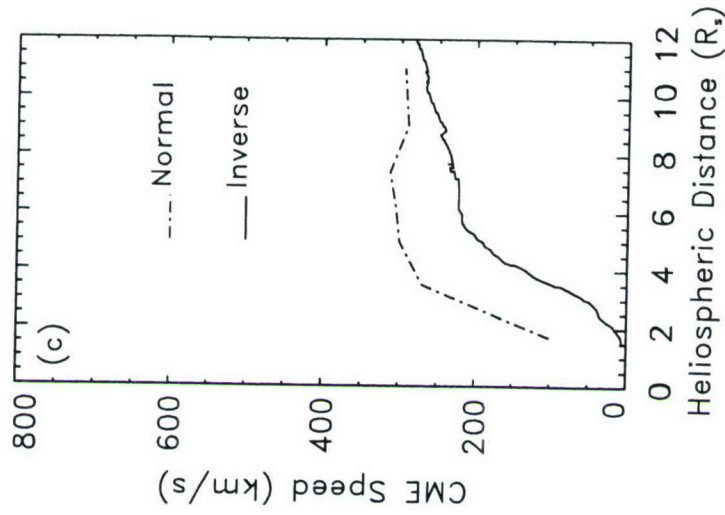
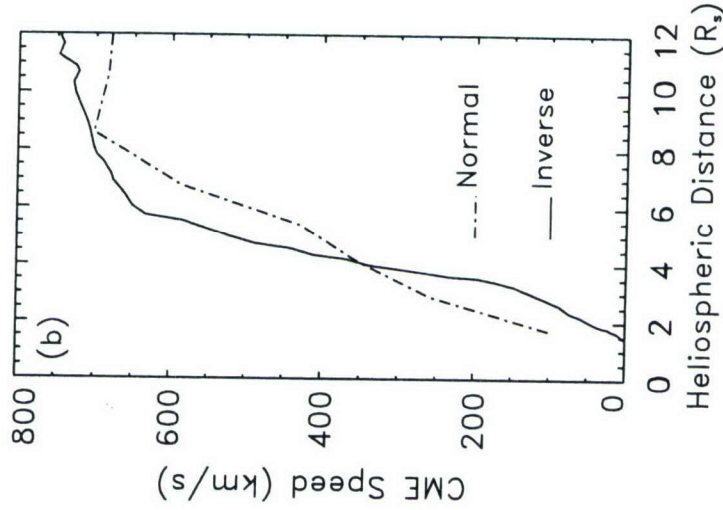
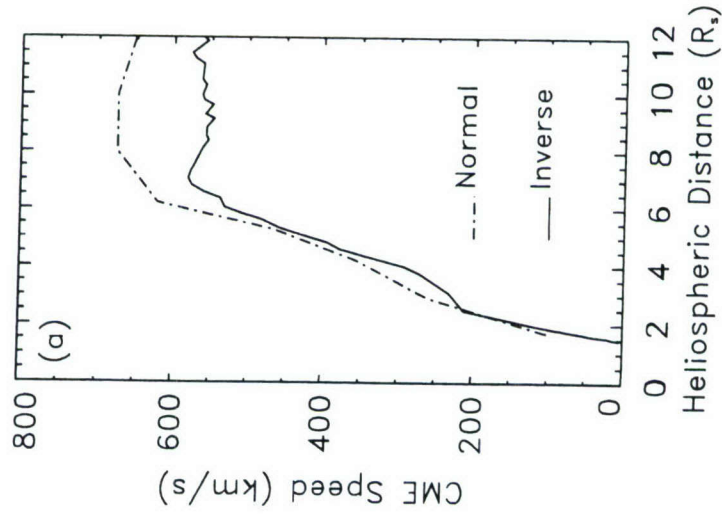
$$T^{n+1} = T^n \left[ 1 + \delta_T \left( 1 - \frac{r^*}{0.85r_f} \right) \right]$$

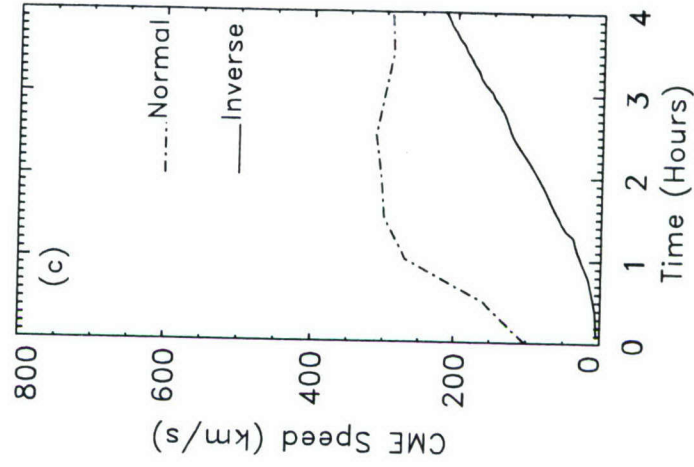
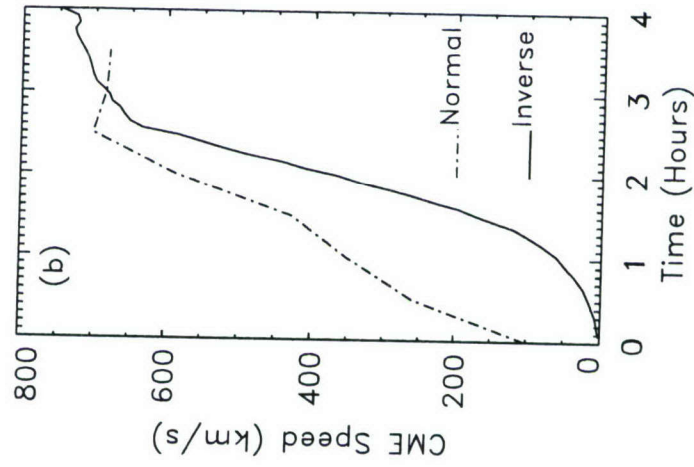
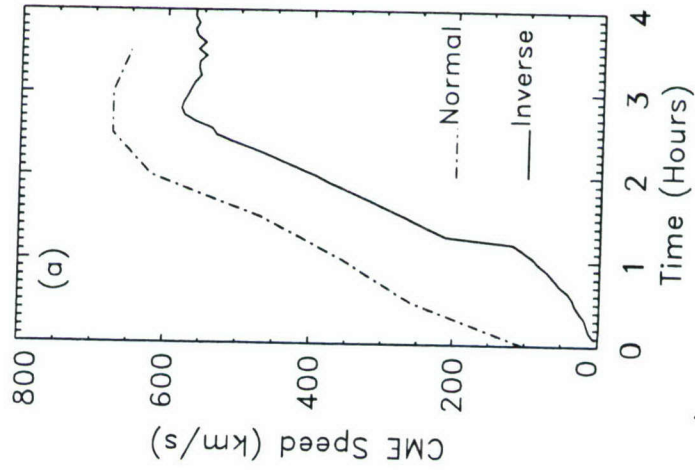
The schematic description of spatial distribution (left) and temporal distribution (right) of perturbed quantities where  $r$  is the flux-rope radius and  $r^*$  is the running variable in the radial direction.

Table II The description of the three perturbations for both **Normal** and **Inverse** flux-rope eruptions of the experiments.

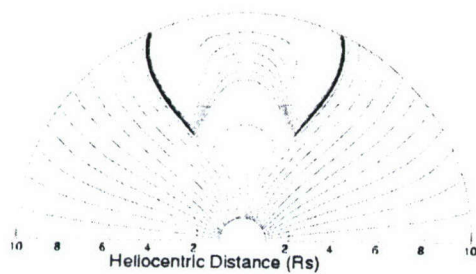
Case	$\delta_B$	$\delta_T$	$\delta_\rho$
1	0.008	0	0
2	0	0.008	0
3	0	0	-0.004



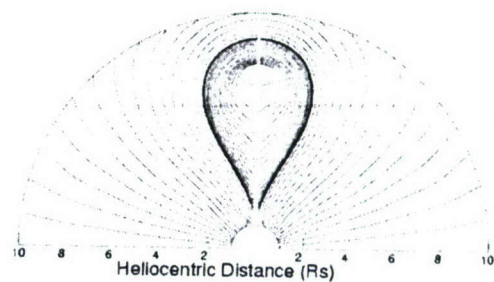




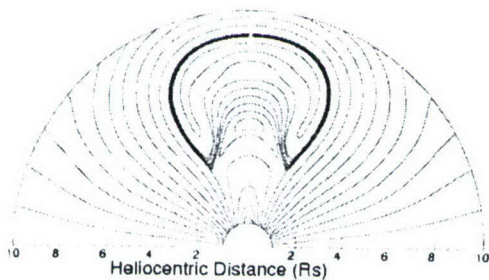
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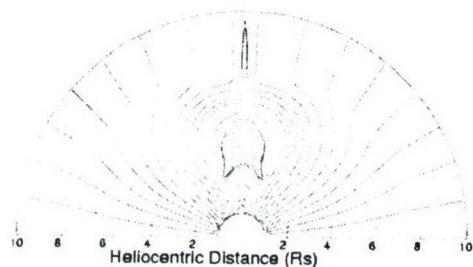
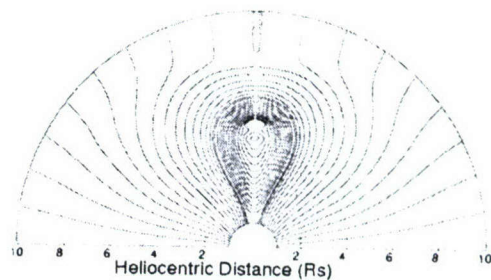
4.0 Hours



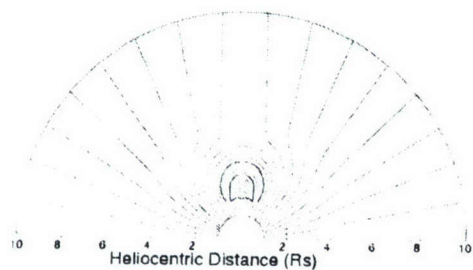
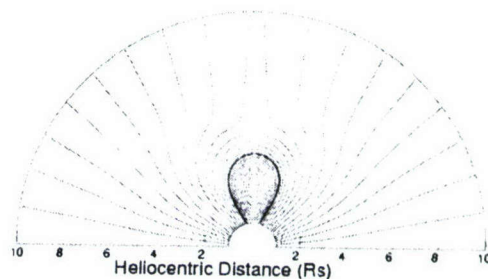
Inverse



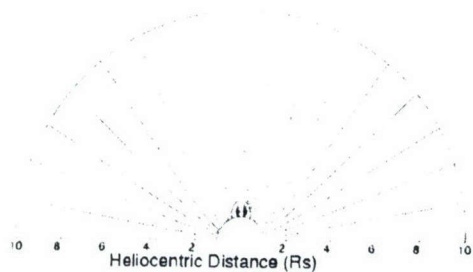
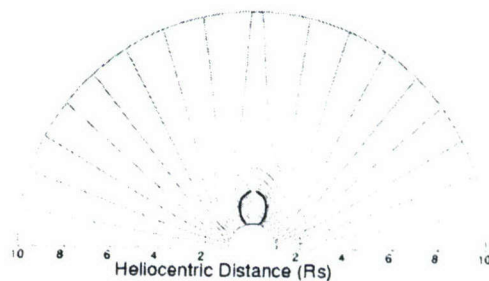
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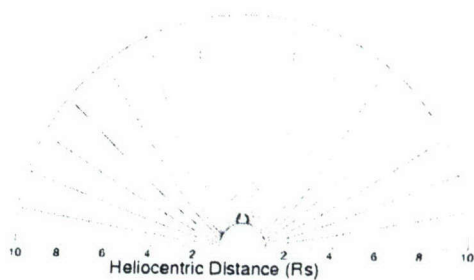
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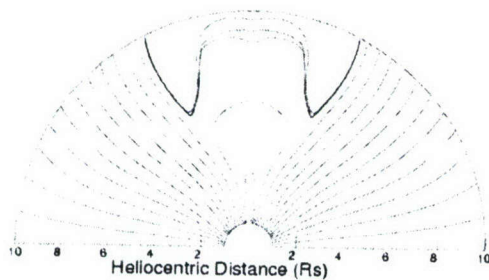


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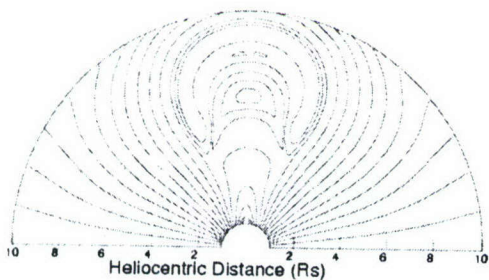
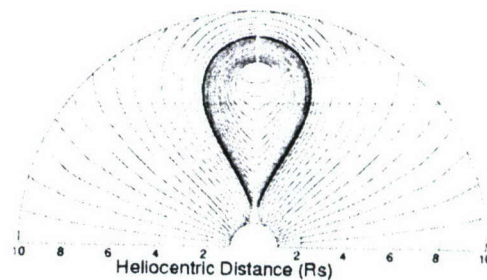


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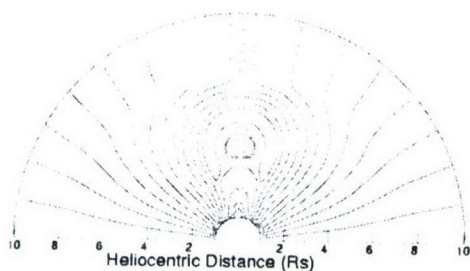
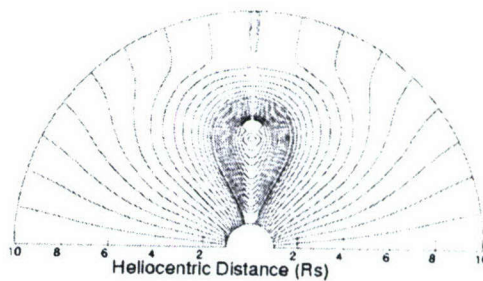


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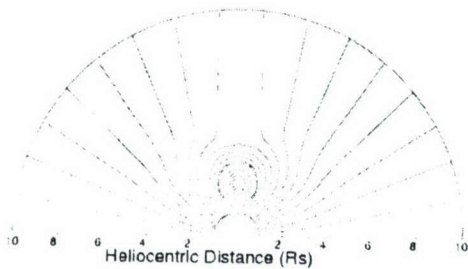
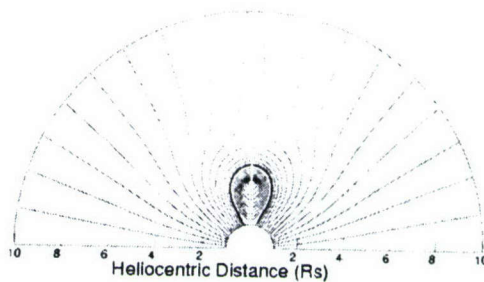
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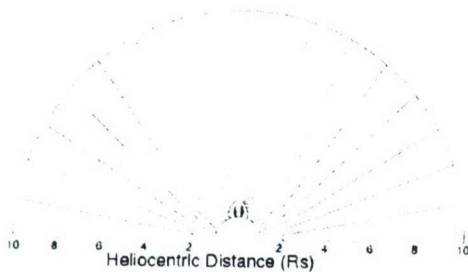
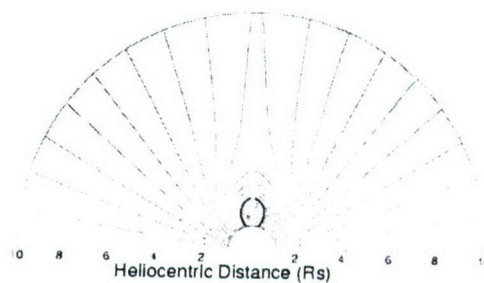
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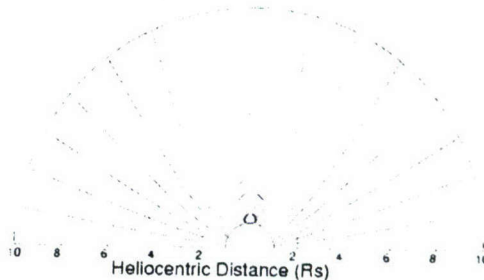
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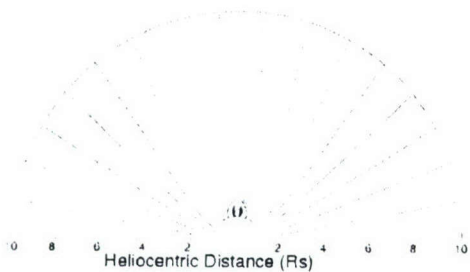
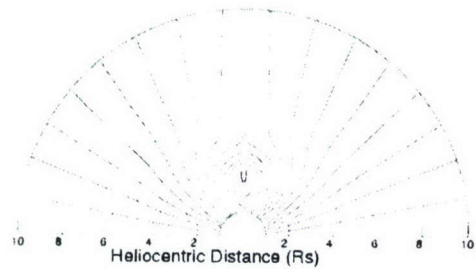
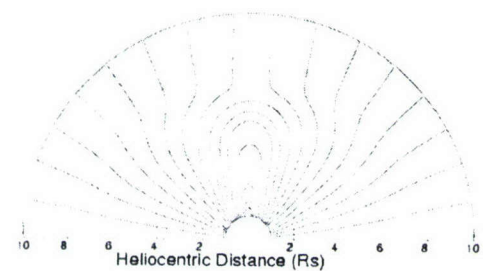
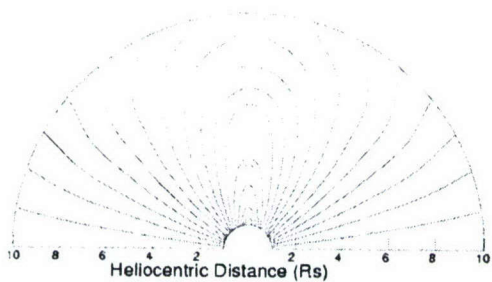
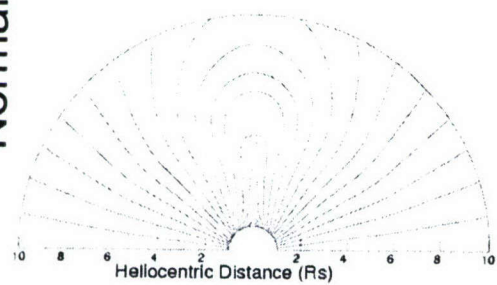
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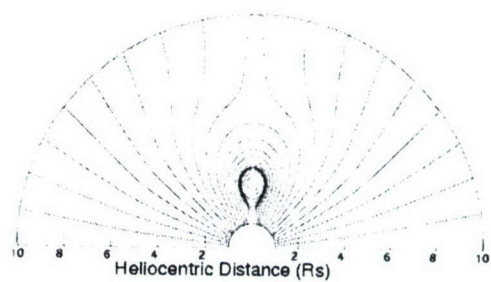
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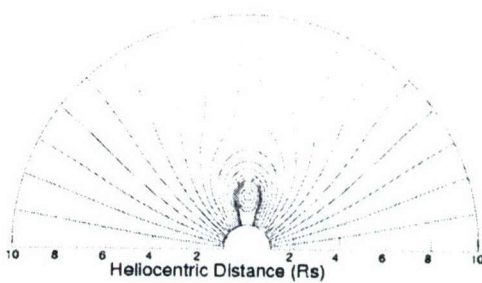
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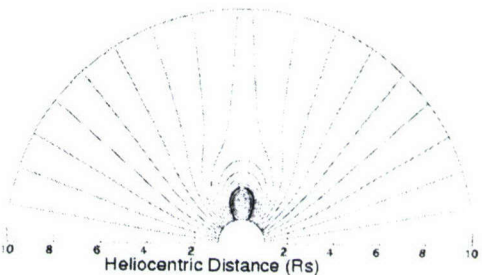
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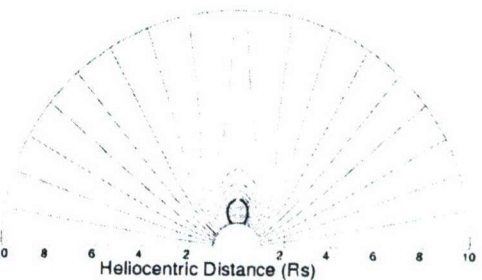
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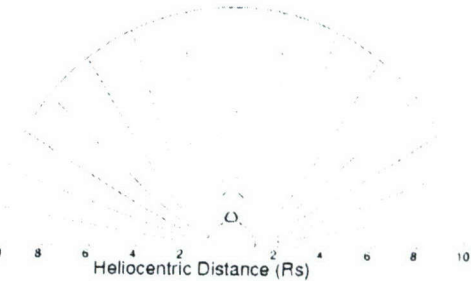
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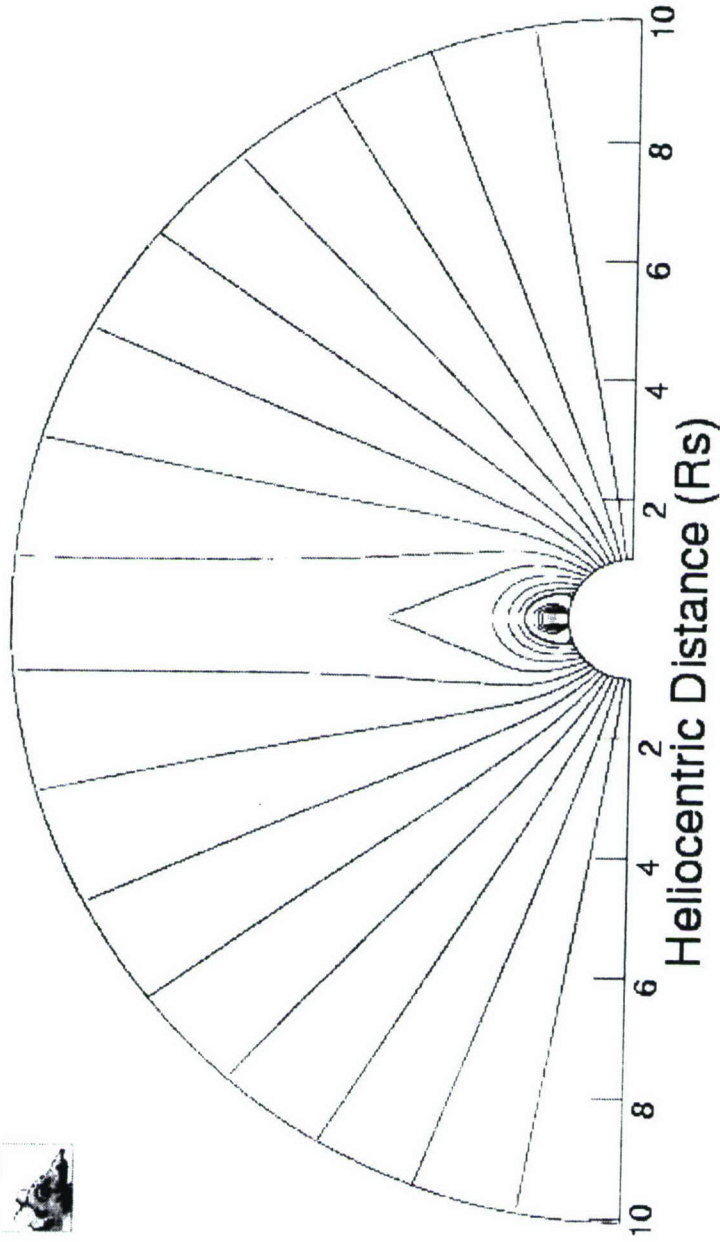


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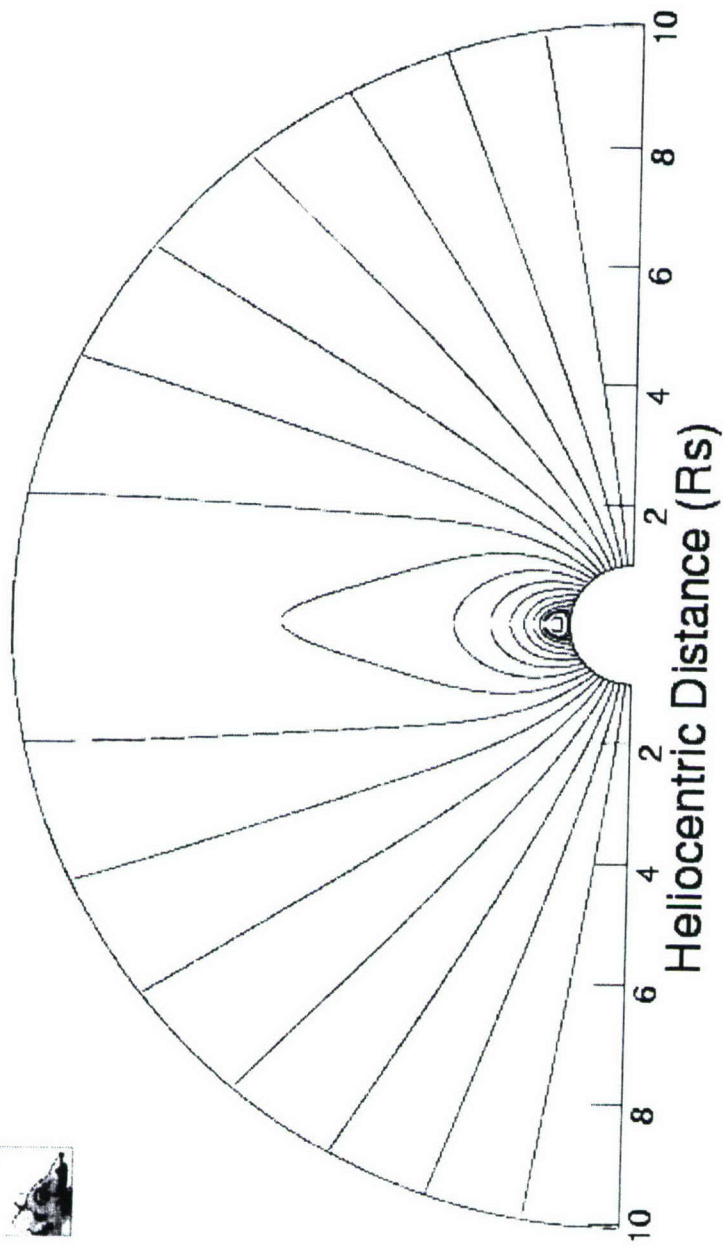
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B-Normal



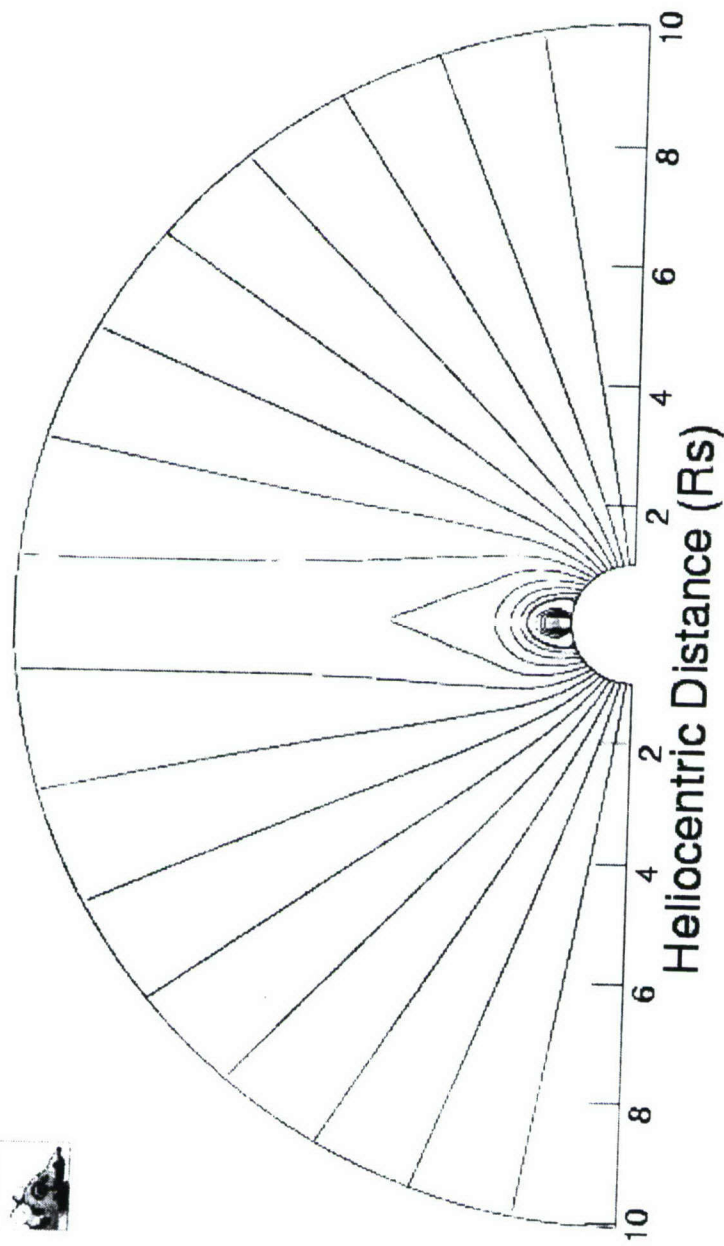


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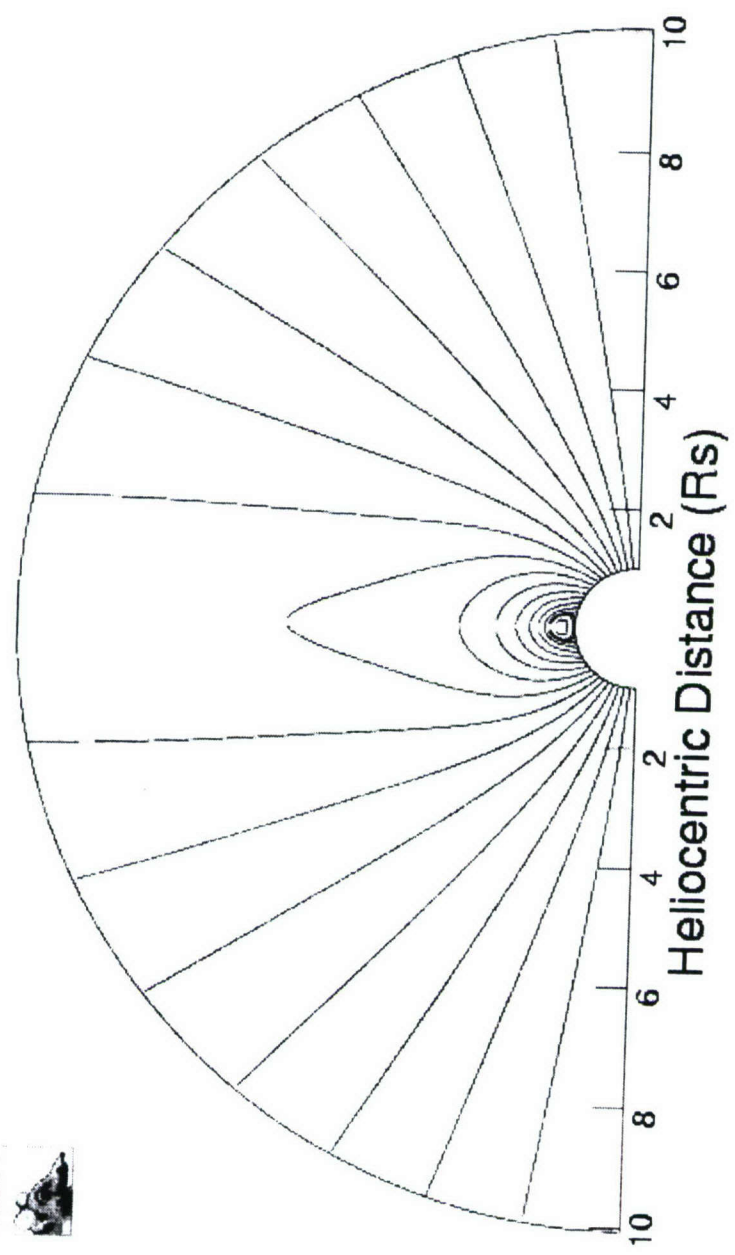




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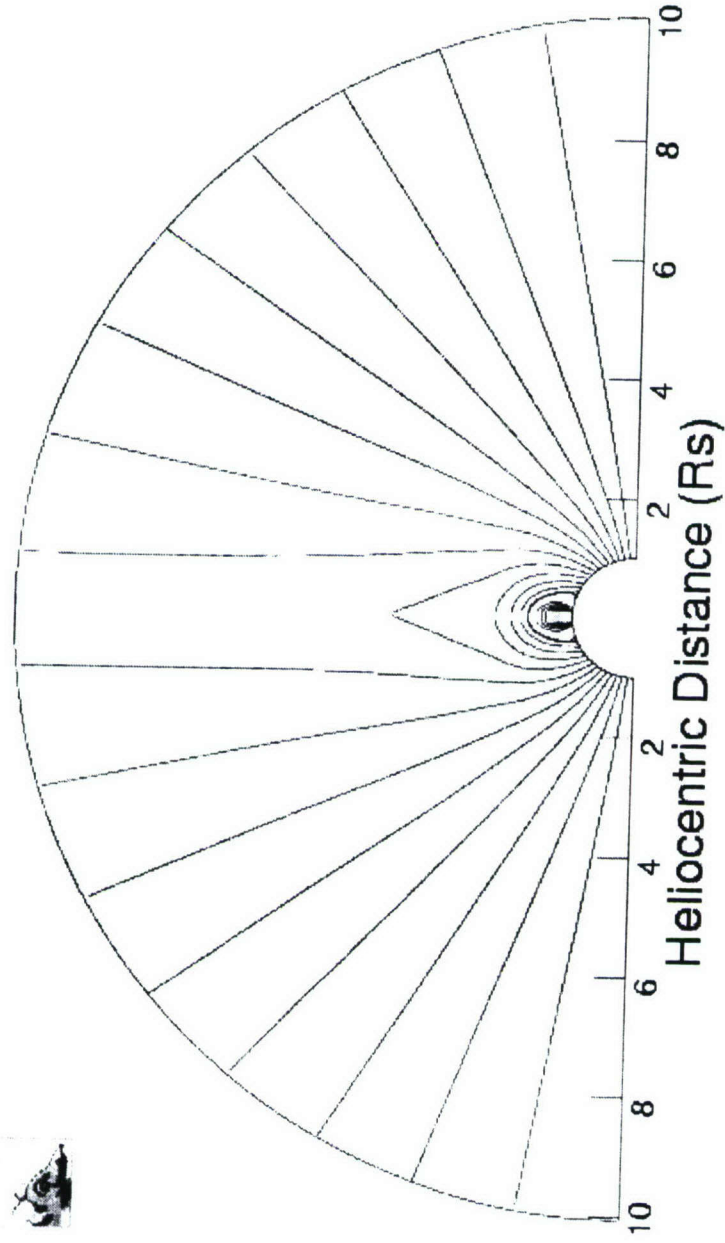


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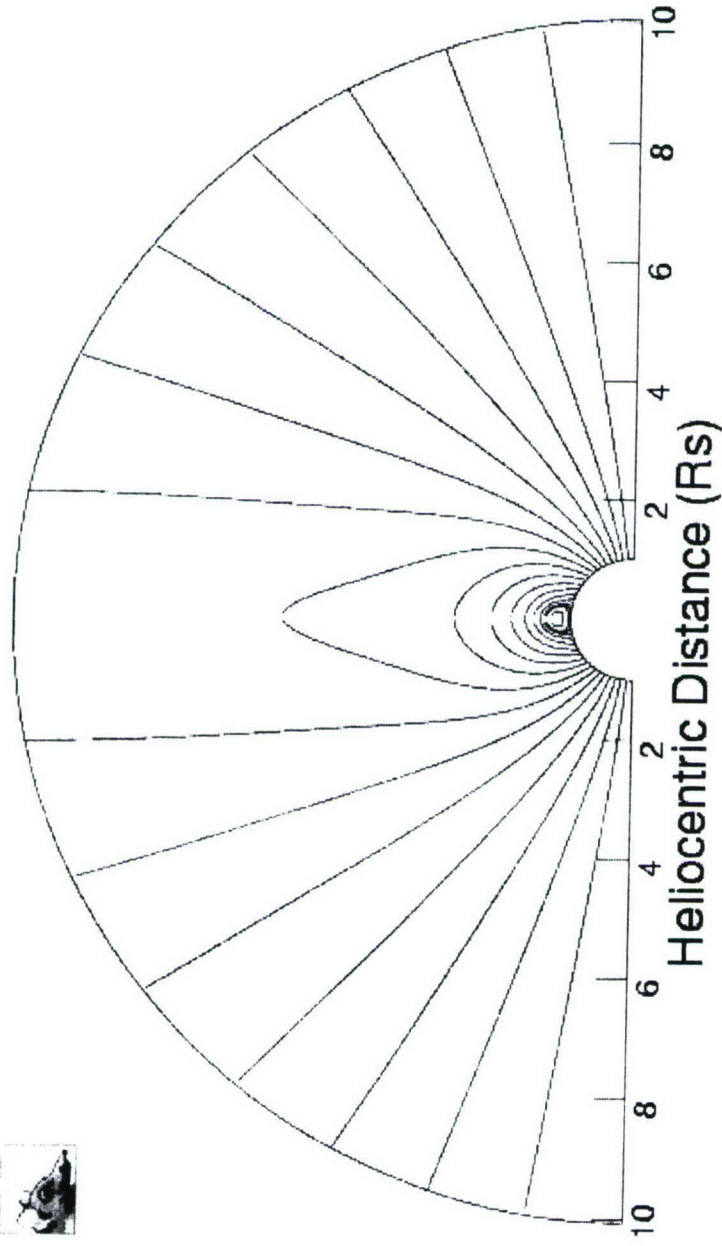




M-Normal



**M-Inverse**



# Summary

In this study, we have employed a 2.5D time-dependent streamer-flux rope MHD model to investigate the role of magnetic reconnection in the CME acceleration processes. To implement this simulation, two distinct types of initial magnetic topology (i.e., Inverse and Normal prominence configuration) are used. The results showed:

- The normal prominence magnetic configuration forms a current sheet in the region of leading edge of the flux-rope, thus the magnetic reconnection occurs to open the closed field of the streamer which allows the flux rope to escape to launch a fast CME as suggested by Low and Zhang (2002). On the other hand, when the initial magnetic topology in the initial prominence configuration, the magnetic reconnection occurs at trailing edge of the flux-rope. In this case, the streamer confined the flux-rope until the magnetic reconnection occurs, this confined force removed which enlarged the CME acceleration (Wu et al. 2004)
- Both inverse and normal prominence configurations are able to produce fast CMEs.
- It is likely that location of the CME induced shock would be close to the Sun's surface for normal prominence configuration than the Inverse one. Because the initial acceleration is higher for the normal configuration.